Extragalactic H I 21-cm absorption line observations with the Five-hundred-meter Aperture Spherical radio Telescope

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1 INTRODUCTION

As the most abundant element in the Universe, hydrogen plays a vital role in the process of stellar and galaxy evolution. With its easily traceable 21-cm hyperfine transition line, the neutral hydrogen (H I) component in galaxies provides a fundamental tool to unveil the phase transition in the interstellar medium (ISM), which is crucial for star formation activities, galaxy kinematics, as well as the large-scale cosmic structures (e.g. see Morganti & Oosterloo 2004; Allison et al. 2016; Curran 2019; and references herein), with samples of $z > 3$ absorbers detected (Uson, Bagri & Cornwell 1991; Kanekar, Chengalur & Lane 2007). Also, such H I absorption lines can bring insights into the continuum sources, usually active galaxy nuclei (e.g. see Maccagni et al. 2017; Morganti & Oosterloo 2018), and references herein, thus revealing possible interactions between star formation and central supermassive black hole activities.

However, only chance alignments between radio-loud active galactic nuclei (AGNs) and foreground or associated neutral hydrogen gas can give rise to H I absorption lines (Morganti & Oosterloo 2018), and sufficiently high column density or large mass of absorbing H I is required for the production of prominent absorption lines (Darling et al. 2011), because of the low Einstein rate of H I emission. Thus, theoretically speaking, due to the occasional nature of H I absorption lines, extensive observations along numerous line of sights are required in order to accumulate a larger sample of H I absorbing systems (Darling et al. 2011). A more practical way to search for H I absorption features is to perform radio observations on damped Lyman-α (DLA) systems, which exhibits high H I column density. However, since the comoving density of low-redshift ($z \leq 0.1$) DLA systems is relatively low (Zwaan et al. 2005), such efforts bear limited results (e.g. see Kanekar et al. 2009).
On the other hand, the absorption line search from AT20G compact radio galaxies using the Australia Telescope Compact Array (ATCA; Allison et al. 2012), the 21-cm Spectral Line Observations of Neutral Gas with the EVLA (21-SPONGE) survey (Murray et al. 2015), the Westerbork Synthesis Radio Telescope (WSRT) radio galaxy surveys (Gerbè, Morganti & Oosterloo 2014; Gerbè et al. 2015; Maccagni et al. 2017), as well as searches by the Australian SKA Pathfinder (ASKAP; Glowacki et al. 2019; Sadler et al. 2020), all targeting known quasars with radio interferometers, have identified dozens of extragalactic H I absorption lines. And similar attempts have been made with the Millennium Arecibo 21-cm Absorption-Line Survey (Heiles & Troland 2003a,b), although using a single-dish telescope. Such targeted observations are mainly focused on extragalactic radio sources with sufficiently high fluxes, thus can only lead to continuum flux-biased H I absorber samples, with galaxies exhibiting weak radio emissions while high H I column densities, which can also give rise to noticeable absorption features, largely missing (e.g. see Sadler et al. 2007).

‘Blind’ searches through H I surveys performed by large aperture single-dish telescopes provide chances to uncover more H I absorption systems. Although such surveys usually utilize the non-tracking, drift scan observing strategy, with integration time for each individual source is limited, the collecting areas of the participating telescopes can still achieve considerable sensitivities. Allison, Sadler & Meekin (2014) identified four H I absorbers, including one previously unknown source, with the archival data of the H I Parkes All-Sky Survey (HIPASS; see Barnes et al. 2001), while Darling et al. (2011), Wu et al. (2015), as well as Song et al. (in preparation) have made attempts to perform ‘blind’ searches for absorption features using part of Arecibo Legacy Fast Arecibo L-Band Feed Array (ALFA) Survey (ALFALFA; see Giovanelli et al. 2005; Haynes et al. 2018) data, respectively, with 10 sources identified in total, including 3 samples remained undetected by other instruments so far, i.e., UGC 00613, CGCG 049-033, and PGC 070403, thus proving the feasibility of searching for new H I absorbers with massive blind sky surveys. Compared with radio interferometers, large single-dish telescopes such as Arecibo can provide better sensitivities and usually higher spectral resolution, which are all crucial to reveal the characteristics of extragalactic H I lines. And although the spatial resolution for such instruments are quite limited compared with interferometers, the chance of having two or more H I absorbing systems lying within the beamwidth with similar redshift is quite low, thus making confusions in source identification unlikely to happen. However, it should be noted that due to temporal variations in spectral baseline commonly seen during drift scans, follow-up observations are needed for reliable characterization of the newly identified absorbers, especially for the weak ones. Also, interferometric mappings are usually required to discern possible fine structures within each absorbing system.

The newly commissioned Five-hundred-meter Aperture Spherical radio Telescope (FAST; Nan 2006; Nan et al. 2011; Jiang et al. 2019) is the largest filled-aperture single-dish radio antenna in the world, with its sensitivity and observable sky coverage both surpassing those of the Arecibo dish. Thus, it is nature to expect better H I absorption observations to be obtained with FAST. Compared with ~several H I absorber discoveries made with existing blind surveys by Wu et al. (2015), Song et al. (in preparation), and Allison et al. (2014), Wu et al. (2015) predicted that the number of absorbing systems detected by FAST should achieve an order of magnitude of at least 10^2 with the upcoming Commensal Radio Astronomy FaST Survey (CRAFTS); see Li et al. 2018), and Yu et al. (2017) gave an expectation of over 1000 extragalactic H I absorption line detections through blind searches in FAST drift scan data. This means that the number of known H I absorbers could be increased by 10 times with FAST, considering currently only ∼100 such systems with redshift z < 1 exist (Chowdhury, Kanekar & Chengalur 2020). And the FAST observations will also complement the ongoing H I absorption surveys performed by the next-generation arrays, including the MeerKAT Absorption Line Survey (MALS; see Gupta et al. 2016), the Widefield ASKAP L-band Legacy All-sky Blind surveyY (WALLABY; see Koribalski et al. 2020), the First Large Absorption Survey in H I (FLASH) by ASKAP (Allison et al. 2020), the Search for H I Absorption with AperTIF (SHARP) survey (Adams et al. 2018), as well as the uGMRT Absorption Line Survey (Gupta et al. 2021), with a more complete Northern sky coverage, better spectral resolution compared with all these surveys, and a similar noise level obtained during much shorter integration times provided by FAST.

In this paper, we report extragalactic H I absorption line observations conducted with FAST on five galaxies identified by Wu et al. (2015) and Song et al. (in preparation), as a pilot study of massive H I absorption line observations by FAST in the near future. The content of this paper is organized as follows. In Section 2, we introduce the galaxy samples, FAST observation settings, as well as the data reduction process in details. The related results are presented and summarized in Section 3. In Section 4, we give discussions, and conclusions are drawn in Section 5. Here, we adopt the ΛCDM cosmology, assuming the cosmological constants as $H_0 = 70\,\text{km\,s}^{-1}\,\text{Mpc}^{-1}$, $\Omega_m = 0.3$, and $\Omega_\Lambda = 0.7$. And in accordance with the ALFALFA survey (Haynes et al. 2018), the optical definition of line speed (δλ/λ) rather than the radio one (δν/ν) is applied throughout our data reduction process.

2 SAMPLE SELECTION AND FAST OBSERVATIONS

2.1 Sample selection

As a pilot study, we tested various observing modes and adopted different backend set-ups to observe the H I spectra during the FAST commissioning phase. The H I absorption systems are particularly suitable for such a purpose, as they contain both continuum sources and spectra lines with know fluxes, thus providing good opportunities for checking telescope pointing accuracy and testing our calibration methods. We select five absorbing systems first identified from 40 per cent of ALFALFA data by Wu et al. (2015), including UGC 00613, ASK 378291.0, CGCG 049-033, J1534+2513, and PGC 070403. Among them, ASK 378291.0 and J1534+2513 were also reported by the WSRT survey for radio galaxies described by Maccagni et al. (2017), with J1534+2513 and PGC 070403 already marked as possible absorptions by Haynes et al. (2011). For each system in our samples, the continuum background is provided by a previously detected radio source, while the absorbing H I gas lies within the source itself. The basic information of these sources are listed in Table 1.

2.2 Observation setting-ups

Four of the five sources have been observed with the tracking mode during the commissioning phase of the FAST telescope. And two of these, ASK 378291.0 and UGC 00613, have been observed on 2017 October 29 and November 5, respectively, with the wide-band receiver covering the 270–1620 MHz frequency range. Since the spectral backend for FAST was still in development at that time, an off-the-shelf N9020A MXA spectrum analyser produced by
Table 1. Basic characteristics of the five targets observed by FAST.

<table>
<thead>
<tr>
<th>Source name</th>
<th>AGC number</th>
<th>RA (J2000)</th>
<th>Dec. (J2000)</th>
<th>$cz$ (km s$^{-1}$)</th>
<th>$S_{1.4GHz, NVSS}$ (mJy)</th>
<th>$S_{1.4GHz, FIRST}$ (mJy beam$^{-1}$)</th>
</tr>
</thead>
<tbody>
<tr>
<td>UGC 00613</td>
<td>613</td>
<td>00$^h$ 59$^m$ 24.62s</td>
<td>+27d 03m 32.6s</td>
<td>13770.07 ± 23.08</td>
<td>112.6 ± 4.0</td>
<td>–</td>
</tr>
<tr>
<td>ASK 378291.0</td>
<td>712.921</td>
<td>10$^h$ 25$^m$ 41.23s</td>
<td>+10d 22m 30.5s</td>
<td>13732.00 ± 26.98</td>
<td>76.6 ± 2.3</td>
<td>73.85 ± 0.145</td>
</tr>
<tr>
<td>CGCG 049–033</td>
<td>250.239</td>
<td>15$^h$ 11$^m$ 31.38</td>
<td>+07d 15m 07.1s</td>
<td>13383.94 ± 35.98</td>
<td>108.8 ± 3.8</td>
<td>79.50 ± 0.151</td>
</tr>
<tr>
<td>J1534+2513</td>
<td>727.172</td>
<td>15$^h$ 34$^m$ 37.62</td>
<td>+25d 13m 11.4s</td>
<td>10180.65 ± 24.98</td>
<td>50.1 ± 1.6</td>
<td>44.58 ± 0.146</td>
</tr>
<tr>
<td>PGC 070403</td>
<td>331.756</td>
<td>23$^h$ 40$^m$ 28.24s</td>
<td>+27d 21m 26.5s</td>
<td>7520.894 ± 118.4</td>
<td>116.6 ± 3.5</td>
<td>–</td>
</tr>
</tbody>
</table>

Notes: *(a) Serial number in the Arecibo General Catalogue (AGC), see Haynes et al. (2018).
(b) Positional coordinates for corresponding optical counterparts retrieved from NASA/IPAC Extragalactic Database (NED, http://ned.ipac.caltech.edu/).
(c) Heliocentric redshift data retrieved from NED.
(d) 1.4 GHz integrated flux data retrieved from the NRAO Very Large Array (VLA) Sky Survey (NVSS) catalogue (Condon et al. 1998, https://www.cv.nrao.edu/nvss).
(e) 1.4 GHz peak flux data retrieved from the VLA Faint Images of the Sky at Twenty-Centimeters (FIRST) survey (Becker, White & Helfand 1995, http://sundog.stsci.edu/cgi-bin/searchfirst). (For comparison only in order to keep consistency with Wu et al. 2015 and Song et al., in preparation; not used in data reduction.)*

2.3 Flux calibration and baseline subtraction

The 19-beam receiver of FAST is equipped with built-in noise diode for signal calibration. As illustrated in Jiang et al. (2020), the diode can be operated at two modes, with characteristic noise temperatures $T_{noise}$ as $\sim 1.2$ and $\sim 12$ K, respectively. For ASK 378291.0, CGCG 049–033, J1534+2513, and PGC 070403 observations, we all adopted the high-temperature mode of the noise diode; and the calibration signal was injected periodically during each observing session. Let OFF be the original instrument reading without noise, ON the reading with noise injected (it should be noted here that throughout this work, we adopt capitalized ON and OFF to describe the noise diode status, while the lower case on/off denote the on/off source position). $T_{noise}$, the pre-determined noise temperature measured with hot loads, which has been shown in Jiang et al. (2020), the system temperature $T_{sys}$ can be computed as

$$T_{sys} = \frac{OFF}{ON - OFF} \times T_{noise}. \quad (1)$$

For each target observed with the 19-beam receiver, we calculated averaged $T_{sys}$ with equation (1) using averages of OFF and ON during each observing session. And as seen in Fig. 2, since the phase of instrumental standing waves changes with each activation of the high-temperature noise, and the wavelength of the standing wave ripples induced by the FAST receiving system can be approximated as $c/2f \sim 1.1$ MHz (where $c$ is the speed of light, and $f \sim 138$ m the focal length of FAST), corresponding to $\sim 140–150$ spectral channels, we performed a several-hundred-channel Gaussian smooth on each set of ON−OFF value, to minimize the phase-shifting effects in calibrations. Then $T_{sys}$ from both polarizations were manually checked, and if they were consistent with each other (which was the case for all four samples observed by the 19-beam receiver), each pair of polarization A and B temperature readings were added together to get the average. Finally, the polarization-averaged $T_{sys}$ was converted to flux densities with pre-measured antenna gain ($\sim 15.7–16.5$ K Jy$^{-1}$).
Table 2. FAST observation setting-ups for the five targets.

<table>
<thead>
<tr>
<th>Source name</th>
<th>Observation mode</th>
<th>Observation session</th>
<th>Receiver</th>
<th>Backend</th>
<th>Band coverage (MHz)</th>
<th>Channel number</th>
<th>Observation date</th>
<th>Duration (s)</th>
</tr>
</thead>
<tbody>
<tr>
<td>UGC 00613</td>
<td>Tracking</td>
<td>1</td>
<td>W</td>
<td>A</td>
<td>1354.3–1360.3</td>
<td>1,001</td>
<td>2017 Nov. 05</td>
<td>1200</td>
</tr>
<tr>
<td>ASK 378291.0</td>
<td>Tracking</td>
<td>1</td>
<td>W</td>
<td>A</td>
<td>1354–1360</td>
<td>1,001</td>
<td>2017 Oct. 29</td>
<td>1200</td>
</tr>
<tr>
<td></td>
<td>Tracking</td>
<td>2</td>
<td>19</td>
<td>S</td>
<td>1000–1500</td>
<td>65,536</td>
<td>2018 Sep. 02</td>
<td>382.5</td>
</tr>
<tr>
<td>CGCG 049-033</td>
<td>Tracking</td>
<td>1</td>
<td>19</td>
<td>S</td>
<td>1000–1500</td>
<td>65,536</td>
<td>2018 Oct. 31</td>
<td>1200</td>
</tr>
<tr>
<td>J1534+2513</td>
<td>Tracking</td>
<td>1</td>
<td>19</td>
<td>S</td>
<td>1000–1500</td>
<td>65,536</td>
<td>2018 Dec. 28</td>
<td>1200</td>
</tr>
<tr>
<td>PGC 070403</td>
<td>Drifting</td>
<td>1</td>
<td>19</td>
<td>S</td>
<td>1000–1500</td>
<td>65,536</td>
<td>2020 Sep. 14</td>
<td>~12</td>
</tr>
</tbody>
</table>

Notes:

a 'W' for the wide-band receiver; '19' for the 19-beam receiver.
b 'A' for Keysight spectrum analyser; 'S' for spectral line backend.

Figure 1. Top panel: The original, unprocessed on-source (blue) and off-source (green) data sample of ASK 378291.0. Bottom panel: Noise diode-calibrated on-source (blue) and off-source (green) data of ASK 378291.0. It can be seen that the flux levels and standing wave behaviours of the two spectra differ from each other significantly, even for the calibrated data. Besides, the off-source data exhibits a higher background level, which is possibly due to temporal instrumental baseline fluctuations, as well as broadband RFI contamination. The red dotted line mark the frequency of the desired absorption feature. Observed frequencies shown in this figure are not heliocentric-corrected.

Figure 2. Averaged on-source spectra of ASK 378291.0 with noise ON (blue, offset in y-axis added for convenience) and OFF (green). The absorption feature can be found near to ~1359 MHz, indicated by the red dotted line. It can be seen that the phase of standing waves imposed on the background continuum flips between the ON and OFF states. Frequencies shown in this figure have not been heliocentric corrected yet.

However, during the ‘early-science’ stage of FAST’s commissioning phase, when the wide-band receiver was still in use, no reliable noise diode was available for flux calibrations. What was worse, the level of instrument reading for N9020A MXA spectrum analyser is often relatively unstable, even within one observing session. Thus, it poses significant difficulties in calibrating the UGC 00613 data. Fortunately, the absorption feature of ASK 378291.0 was detected with both the wide-band receiver and the 19-beam array. Thus, we adopted the measured flux density of ASK 378291.0 from the latter instrument as reference to re-scale and calibrate the wide-band receiver observations. And since only one polarization was available during this period, we just assumed that the single set of flux densities applies for both polarizations. Although such operations could bring considerable compromise to the accuracy of our final products, it is the only sensible method that could produce reasonable results generally compatible with the ALFALFA observations, as provided by Wu et al. (2015) and Song et al. (in preparation). More details of our results can be found in Section 3.1.1.
Once the flux densities were calculated via calibration, the observed frequencies were Doppler-corrected to the heliocentric frame. Since the fully automatic data reduction pipeline for FAST H1 observations (Zhang et al. 2019) is still in development, all baseline subtraction and RFI flagging works involved in this paper were performed manually. For every source, the overall shape of the background continuum (‘baseline’) curve was determined by calculating the median value for every channel during each entire session, similar to the process utilized by H1 Parkes All Sky Survey (HIPASS, see Barnes et al. 2001). As for the baseline reading for the frequency channels in which the desired H1 absorption lines resides, a cubic spline interpolation was performed. Once the baseline was subtracted, strong RFI was marked by visual inspections. Finally, one- or two-component Gaussian fitting was applied to data not contaminated by RFI, depending on the line profiles, to get the spectral parameters. Here, considering the baseline instabilities between on- and off-source positions as described in Subsection 2.2, continuum flux densities measured by the NVSS survey (rather than direct measurements by FAST) were adopted for the calculation of optical depth and other line characteristics of each H1 absorbing feature.

Here, it should be noted that since the NVSS was performed more than two decades ago, it is quite possible that the continuum flux level for the observed sources varies during this period. In fact, as shown in Table 1, the continuum flux of CGCG 049-03 does show a difference as large as 27 per cent between NVSS and the earlier VLA Faint Images of the Sky at Twenty-Centimeters (FIRST) survey (Becker et al. 1995) results, which may be due to possible source variability, or the existence of the extended jet in this source (Bagchi et al. 2007), combined with the varied VLA configurations adopted by different surveys. (While for ASK 378291.0 and J1534+2513, which have also been observed by both NVSS and FIRST, measurements obtained by these two surveys are generally compatible, with flux variations less than \( \sim 10\) per cent.) Also, the beam size difference between NVSS survey and FAST, along with the unfilled \( u-v \) coverage of VLA may cause the NVSS flux to be somewhat underestimated for this work. Thus, as mentioned in Section 2.1, our strategy to calculate absorption parameters (including H1 column density \( N_{\text{H1}} \) and optical depth \( \tau \)) based on existing survey results is only a compromise to the not-so-stable behaviours of our newly commissioned instrument, and is adopted in reference to the solution utilized by various single-dish absorption line observations, such as Darling et al. (2011), Grasha et al. (2019), as well as Zheng et al. (2020).

3 OBSERVING RESULTS

We successfully detected H1 absorption features in all of the five sources firstly identified with 40 per cent of the ALFALFA data by Wu et al. (2015) and Song et al. (in preparation). The properties of each detected source estimated using FAST are summarized in Table 3. The column densities of H1 towards absorption line \( N_{\text{H1}} \) in this table are calculated with the following equation

\[
N_{\text{H1}} = 1.823 \times 10^{14} \frac{T_{\text{P}}}{f} \int \tau \, dv \, \text{cm}^{-2}.
\]  

Here, \( T_{\text{P}} \) denotes the H1 spin temperature, \( f \) the covering factor of the radio continuum source, with a typical value of unity assumed (e.g. see Maccagni et al. 2017). Since no source observed by this work is associated with known DLA system (Wu et al. 2015), which usually exhibits a higher \( T_{\text{P}} \) in the order of \( \sim 10^3 \) K (e.g. see Chengalur & Kanekar 2000; Darling et al. 2011; Curran 2019), a typical value of \( T_{\text{P}} \sim 100 \) K is assumed. And the optical depth \( \tau \) of the H1 absorber can be calculated as

\[
\tau = -\ln \left( 1 + \frac{S_{\text{H1}}}{S_{1400\,\text{MHz}}} \right).
\]  

where \( S_{\text{H1}} \) is the depth of the absorption line (in negative value) and \( S_{1400\,\text{MHz}} \), the 1400 MHz flux of the continuum source as provided by NVSS survey. All errors listed in Table 3 have been evaluated with the spectral resolution and root mean square (rms) level of each observing session, the calibration accuracy of the 19-beam receiving system (\( \sigma_{\text{cal}} \sim 0.8\text{--}1.8 \) per cent, as noted by Jiang et al. 2020), the \( \sigma_{\text{v,4 GHz}} \), value of NVSS or VLA FIRST measurements, as well as the line profiles, in reference to the method adopted by Koribalski et al. (2004), as follows

\[
\sigma_{\text{cal}} = 3 \sqrt{\frac{\sigma^2_{\text{v,4 GHz}}}{N}} \times \frac{S}{N}.
\]  

3.1 Note on individual sources and comparisons with Arecibo/WSRT observations

3.1.1 ASK 378291.0

The ASK 378291.0 at redshift \( z = 0.04580 \) has been classified as a newly defined Fanaroff–Riley type 0 (Baldi, Capetti & Massaro 2018) galaxy by Cheng & An (2018), with a matched NVSS source exhibiting a \( \sim 76.6 \) mJy flux at the 1.4-GHz band, and a well-aligned jet and counterjet (Cheng & An 2018) on its position. Its absorption line was detected by both of the wide-band receiver, as well as the 19-beam array of FAST. As shown in Fig. 3, the frequency and the double-horned profile in data from FAST and Arecibo show rough agreements to each other, thus proving the robustness of FAST detections. However, the low frequency peak in FAST data shows a significantly larger depth than ALFALFA’s results. As can be seen in Fig. 4, this discrepancy exist even if the FAST data are binned to resolution similar to that of ALFALFA. Considering the fact that spectra from both of the 19-beam and the wide-band receiver of FAST unveiled similar line profiles, such a discrepancy between data acquired by the two telescopes could be due to footprint positions.
Table 3. FAST observations of the five extragalactic H I absorbers firstly identified with ALFALFA data. All parameters are listed in the rest frame of the absorbers. All background rms data are measured with the original velocity bin. Parameters including $c_{\text{peak}}$, FWHM, $S_HI,\text{peak}$, $\tau_{\text{peak}}$, and $\int \tau dv$ are calculated with the fitted line profiles.

<table>
<thead>
<tr>
<th>Source name</th>
<th>rms (mJy)</th>
<th>Gaussian component</th>
<th>$c_{\text{peak}}$ (km s$^{-1}$)</th>
<th>FWHM (km s$^{-1}$)</th>
<th>$S_{HI,\text{peak}}$ (mJy)</th>
<th>$\tau_{\text{peak}}$</th>
<th>$\int \tau dv$ (km s$^{-1}$)</th>
<th>$N_{HI}$ (10$^9$ cm$^{-2}$)</th>
</tr>
</thead>
<tbody>
<tr>
<td>UGC 00613</td>
<td>6.49</td>
<td>A</td>
<td>13 956.5 ± 0.9</td>
<td>26.36 ± 1.74</td>
<td>−64.28 ± 6.61</td>
<td>0.801 ± 0.145</td>
<td>21.24 ± 4.10</td>
<td>38.72 ± 7.95(T/100 K)</td>
</tr>
<tr>
<td>ASK 378291.0</td>
<td>0.89</td>
<td>A</td>
<td>13 672.3 ± 0.2</td>
<td>14.06 ± 0.36</td>
<td>−42.24 ± 1.22</td>
<td>0.607 ± 0.051</td>
<td>9.059 ± 0.753</td>
<td>16.58 ± 1.46(T/100 K)</td>
</tr>
<tr>
<td></td>
<td></td>
<td>B</td>
<td>13 702.2 ± 0.1</td>
<td>14.09 ± 0.30</td>
<td>−45.93 ± 1.28</td>
<td>0.683 ± 0.061</td>
<td>9.724 ± 0.887</td>
<td>17.72 ± 1.72(T/100 K)</td>
</tr>
<tr>
<td></td>
<td></td>
<td>Total</td>
<td>13 687.3 ± 0.1</td>
<td>44.00 ± 0.30</td>
<td></td>
<td></td>
<td>18.83 ± 1.16</td>
<td>34.33 ± 2.26(T/100 K)</td>
</tr>
<tr>
<td>CGCG 049–033</td>
<td>0.45</td>
<td>A</td>
<td>13 454.4 ± 1.0</td>
<td>52.71 ± 2.07</td>
<td>−5.851 ± 0.468</td>
<td>0.055 ± 0.005</td>
<td>3.069 ± 0.285</td>
<td>5.595 ± 0.553(T/100 K)</td>
</tr>
<tr>
<td>J1534+2513</td>
<td>0.37</td>
<td>A</td>
<td>10 089.0 ± 0.1</td>
<td>5.148 ± 0.240</td>
<td>−14.97 ± 0.487</td>
<td>0.355 ± 0.019</td>
<td>2.291 ± 0.130</td>
<td>4.176 ± 0.253(T/100 K)</td>
</tr>
<tr>
<td></td>
<td></td>
<td>B</td>
<td>10 107.9 ± 0.2</td>
<td>13.76 ± 0.33</td>
<td>−15.32 ± 0.487</td>
<td>0.365 ± 0.020</td>
<td>4.713 ± 0.296</td>
<td>8.593 ± 0.575(T/100 K)</td>
</tr>
<tr>
<td></td>
<td></td>
<td>Total</td>
<td>10 098.5 ± 0.1</td>
<td>27.53 ± 0.29</td>
<td></td>
<td></td>
<td>7.011 ± 0.324</td>
<td>12.78 ± 0.63(T/100 K)</td>
</tr>
<tr>
<td>PGC 070403</td>
<td>2.53</td>
<td>A</td>
<td>7694.2 ± 4.2</td>
<td>74.54 ± 8.37</td>
<td>−9.602 ± 2.539</td>
<td>0.086 ± 0.024</td>
<td>6.643 ± 2.314</td>
<td>12.11 ± 4.49(T/100 K)</td>
</tr>
</tbody>
</table>

Figure 3. Top panel: H I absorption feature ASK 378291.0. The black line shows observations taken by the FAST 19-beam receiver with its spectral line backend on 2018 September 2, while the green curve denotes data taken from the ALFALFA survey. The magenta dotted line is the fitting result by dual-Gaussian functions, while the cyan curve shows the fitting residual with a 15 mJy offset for clarity. The regions shaded with light yellow are manually flagged RFI, which have been all excluded in the fitting procedure and rms computations. The red dotted line shows the optical redshift, while the blue line marks the zero flux level. Bottom panel: The same target observed by the wide-band receiver along with the N9020A MXA spectrum analyser on 2017 October 29, with the original, uncalibrated instrumental readings measured in mW shown along the y-axis. This set of data marks the first successful detection of extragalactic H I absorption line by the FAST telescope.

Figure 4. Comparison between FAST 19-beam, WSRT, and ALFALFA observations of ASK 378291.0. The black line shows data taken by FAST, while the green–grey line is observations from WSRT, as shown in Maccagni et al. 2017, and the green line the ALFALFA data. The dotted cyan line denotes binned FAST data, with a final resolution similar to WSRT’s 16 km s$^{-1}$, while the dotted magenta line denotes the four-channel binned FAST data to achieve a spectral resolution comparable with that of ALFALFA. The red line denotes the optical redshift and the blue line marks the zero level.

in the ALFALFA survey, since the source might not pass through the right center of the receiver beam during those drift scans, thus leading to off-axis measurements, as well as compromised flux and line profile estimations. Besides, the same peak in the 19-beam data shows a clear spike at 1358.3 MHz, which cannot be identified in the spectrum taken by the wide-band receiver. Since multiple RFI features exist within the desired frequency range, it is reasonable to assume that such a spike is due to undesired interference overlaid on the double-horned absorption line. Thus, this spike was excluded in the fitting procedure of the line shape, as well as optical depth calculation.

Compared with the results listed in Song et al. (in preparation), which adopted a single-Gaussian function for line profile fitting, and yielding a line depth of $S_{HI,\text{peak}} \approx -23.5$ mJy, an FWHM value of $\sim 56.3$ km s$^{-1}$, and $\int \tau dv \approx 21.72$ km s$^{-1}$, the FAST observations of ASK 378291.0 show a slightly smaller integrated optical depth, along with a narrower FWHM. However, it should be noted that the ALFALFA fitting result is based on the continuum flux provided by the VLA Faint Images of the Sky at Twenty-Centimeters...
It should also be noted that this absorption feature has also been detected by the WSRT absorption line survey, with an FWHM of $\sim 66.95 \text{ km s}^{-1}$, and $\int \tau d\nu \approx 9.62 \text{ km s}^{-1}$. Maccagni et al. (2017). Fig. 4 shows a comparison between the FAST and WSRT data. It can be seen that the WSRT data only exhibit a single-peaked structure, rather than a double-horned one as shown by FAST and ALFALFA. And the maximum absorption measured by WSRT is only $-13.55 \text{ mJy}$, which is nearly 3.4 times less than measurements performed by FAST. When channel-binned to a velocity resolution equivalent to WSRT, a single-peak absorption line emerges in the FAST data, although with a lopsided profile, and a deeper feature of $-25.08 \text{ mJy}$. Thus, even adopting the higher continuum flux value from Maccagni et al. (2017) ($S_{1400MHz} \sim 92.83 \text{ mJy}$) compared with NVSS or VLA FIRST catalogues, the FAST data can still lead to an $\int \tau d\nu \approx 17.22 \text{ km s}^{-1}$, which is nearly 1.8 times higher than the WSRT integrated optical depth. Such a clear difference could be attributed to the incomplete $u-v$ coverage of a radio telescope array, which may result in compromised flux measurements.

3.1.2 UGC 00613

The absorption feature in UGC 00613, a flat spectrum radio galaxy with extremely diffuse envelope and faint lobes at $z = 0.04593$ (Condon, Frayer & Broderick 1991), was first reported by Wu et al. (2015) and Song et al. (in preparation), with a line depth of $S_{\text{HI, peak}} \sim -65.0 \text{ mJy}$, an FWHM of $\sim 33.6 \text{ km s}^{-1}$, and $\int \tau d\nu \approx 31.14 \text{ km s}^{-1}$. The FAST telescope performed tracking observations on this source with its early-science phase, with its wide-band receiver. Due to the lack of reliable noise diode in this period, the flux of UGC 00613 obtained by the N9020A MXA spectrum analyser were estimated in reference to observations of ASK 378291.0 shown in the upper panel of Fig. 3. Firstly, the baseline-subtracted wide-band-receiver data of ASK 378291.0 (lower panel of Fig. 3) was re-scaled according to the flux level of the same source measured by the 19-beam receiver; then, the same scaling factor for conversion from N9020A MXA’s instrumental reading to flux density (measured in mJy) was applied to UGC 00613 observations. Although such a calibration process is not strictly accurate, the resulted line depth, $\sim -64.34 \text{ mJy}$, is in good agreements with the Arecibo data. The discrepancy between line width estimations of the two telescopes may arise from the higher spectral resolution of FAST, which could lead to a better estimation of the line profile; as well as the relatively high rms level in the N9020A MXA data ($\sim -6.49 \text{ mJy}$, which is an order-of-magnitude higher than that of the 19-beam data), which brings more uncertainties in identification of the wing structure, as can be seen from Fig. 5.

It should be noted that as can be seen in Fig. 5, and also noted by Song et al. (in preparation), the centre of the UGC 00613 absorption feature is redshifted from the optical redshift by as much as $\sim 186 \text{ km s}^{-1}$ ($\sim 188 \text{ km s}^{-1}$ for Arecibo data). According to Wegner et al. (1999), although the SNR of the optical data used for $cz$ estimation is not high enough to accurately measure the velocity dispersion of the spectral lines, the redshift of this galaxy can be reliably determined as $cz = 13770 \pm 23 \text{ km s}^{-1}$. Thus, it is unlikely that such a large offset (which is at $> 8 \sigma$ level of the optical measurement) between H I line and optical velocities arise from observation errors. Rather, this discrepancy may indicate the existence of unsettled infalling gas cloud or other similar structures around this AGN (e.g. see Maccagni et al. 2017 and references herein), and further interferometric spectral line observation is required to reveal the details of the absorbing gas in this galaxy.

3.1.3 CGCG 049-033

The $z = 0.04464$ elliptical galaxy CGCG 049-033 in Abell 2040 cluster possesses a highly asymmetric Fanaroff-Riley type II AGN, comprising of one of the largest radio jet yet discovered, a $>10^9 M_{\odot}$ central black hole, and intense polarised synchrotron radiation extending to a distance of $\sim 440 \text{ kpc}$ away from the galactic centre (Bagchi et al. 2007). The H I absorption feature with a maximum line depth $S_{\text{HI, peak}} \sim -4.3 \text{ mJy}$, an FWHM $\sim 123.9 \text{ km s}^{-1}$, and $\int \tau d\nu \approx 7.23 \text{ km s}^{-1}$ reported by Song et al. (in preparation) is consistent with the general trend that early-type galaxies are in lack of abundant gaseous contents.

The FAST telescope conducted targeted observations twice towards CGCG 049-033, on 2018 October 31st and December 28th, respectively, each was comprised of 1200 s of on source exposure time. Although weak, the expected absorption feature did appear at the right frequency channels in both set of data, thus confirming the correctness of our detection. However, since the December observations suffered stronger RFI contamination, they have not been invoked in our analysis to get the line profile, flux, optical depth, as well as other characteristics of this H I absorber.

As it can bee seen in Fig. 6, the H I absorption line from CGCG 049-033 is the weakest among all five sources detected by Arecibo and confirmed by FAST, with a maximum depth of 5.851 mJy only, as measured by FAST. However, the integrated optical depth, 3.069 km s$^{-1}$, is less than half of that of the ALFALFA results. Even if the lower continuum flux provided by VLA FIRST, as adopted by Song et al. (in preparation), is applied to FAST data, the resulting $\int \tau d\nu$ is still smaller than ALFALFA. This discrepancy should be the result of the much narrower line structure observed by FAST, and the spike-like RFI right on the low-frequency part of the absorption line further affects the fitting of the line profile, resulting in a FWHM...
less than half of Arecibo’s results. Since the SNR for this source in ALFALFA data is quite low, considering the rms level $\sim 2.2$ mJy is almost the half of the peak estimation by Song et al. (in preparation), the ALFALFA line profile itself exhibit considerable uncertainties. Thus, the low-noise FAST data should be considered more reliable, thanks to the better sensitivity and spectral resolution of the FAST telescope.

3.1.4 J1534+2513

The existence of H\textsc{i} absorption feature in radio source J1534+2513 at a redshift of $z = 0.03396$ was first proposed by Haynes et al. (2011), and later confirmed by Wu et al. (2015) and Song et al. (in preparation), with a line depth of $S_{\text{H}i, \text{peak}} \sim -23.0$ mJy, FWHM $\sim 26.4$ km s$^{-1}$, and $\int r d\nu \approx 18.31$ km s$^{-1}$. This absorber has also been detected by the WSRT survey, with a wider line width (FWHM $\sim 40.88$ km s$^{-1}$), and an integrated optical depth of 10.92 km s$^{-1}$, against a 43.39 mJy background (Maccagni et al. 2017).

As can be seen in Fig. 7, the most prominent feature clearly shown in the FAST spectrum should be the double-horned structure over a relatively narrow $\sim 27.53$ km s$^{-1}$ frequency range, which is only barely seen with the $\sim 5.3$ km s$^{-1}$ resolution of the ALFALFA data (Song et al., in preparation), and the $\sim 16$ km s$^{-1}$-resolution WSRT survey (Maccagni et al. 2017, see Fig. 8), thus demonstrating the necessity of performing observations with finer spectral resolution. And the centroid of this absorption line is again significantly deviates from the optical redshift, although in this case, the radio line is blueshifted by $82.15$ km s$^{-1}$ compared with the optical $z$, although the extent of discrepancy is smaller than that of UGC 00613. Considering the error in $cz$ measurement for J1534+2513 is as small as $<2.7$ km s$^{-1}$ (which is the smallest among all five sources discussed in this work), such a line velocity difference must be originated from the absorbing H\textsc{i} content itself. Again, this may indicate the existence of H\textsc{i} outflows along the line-of-sight and interferometer observations should be of help to unveil the nature of the absorbing gas.

Compared with Arecibo and WSRT data, the spectral parameters of J1534+2513 obtained with FAST exhibit a narrower width, a shallower peak depth, as well as a smaller optical depth, compared with the other two sets of data. And although the FAST line profile is significantly deeper than WSRT’s $<10$ mJy (which still holds true if these data are binned to WSRT resolution), the line depth taken by ALFALFA survey is even larger. Besides, the distance between the two spectral peaks in the WSRT data is much larger than FAST and Arecibo’s results. In one sentence, the existing three data sets do not coincide with each other, even when taken spectral resolution into considerations. A possible explanation for such an inconsistency between different instruments could be the errors induced by a relatively high rms level compared with line depth for ALFALFA (which has a baseline fluctuation as large as nearly 10 mJy for frequencies around the J1534+2513 absorption feature, as illustrated in Figs 7 and 8) and WSRT, which can make their results less accurate. Also, the interferometric nature of WSRT observations could bring extra compromise to flux estimations.

3.1.5 PGC 070403

It was Haynes et al. (2011) who put forward the first indication on the existence of $z = 0.0251$ absorption feature PGC 070403, and later the same structure has been tentatively identified by Wu et al. (2015) and Song et al. (in preparation) from the $\alpha.40$ data, with an estimation of line depth as $S_{\text{H}i, \text{peak}} \sim -12.9$ mJy, FWHM $\sim 88.9$ km s$^{-1}$, and $\int r d\nu \approx 11.11$ km s$^{-1}$. As noted by Song et al. (in
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Figure 9. Drift scan data of PGC 070403 from both polarizations. The y-axis is shown in the original instrumental reading. The black line shows the average of ∼12 s original data taken around the transit time of the targeted source, Gaussian-smoothed by four spectral channels to highlight possible line features. The blue and green lines are four-channel Gaussian-smoothed, ∼12 s-averaged data taken before and after source transit, respectively, with ∼10 per cent of adjustments on the y-axis readings for clarity. It can be seen that a ‘dip’ appears at ∼1384.9 MHz around the transit time on both polarizations. And the red dotted line denotes the \( c_\text{peak} \) of the expected absorption feature as measured by Arecibo, which is almost coincident with the position of the transit ‘dip’.

preparation), similar to the case of UGC 00613, this absorption line is also redshifted from the optical \( c_\text{z} \) by as many as ∼191.5 km s\(^{-1}\). However, it is also noticed that the error of optical \( c_\text{z} \) measurement is as large as 118.4 km s\(^{-1}\), which means that such a large discrepancy may be due to observational uncertainties.

Among all of the five H I absorbing sources observed by the FAST telescope and mentioned in this paper, PGC 070403 is the one with the second-weakest line depth, and the only sample observed with the drift scan mode by the FAST telescope. With a beamwidth of ∼2.9 arcmin at ∼1400 MHz band, the central feed of the 19-beam receiver can observe each source for ∼12 s during a one-pass scan. As can be seen in Fig. 9, the absorption ‘dip’ only appears during the ∼12 s period around the transit time of PGC 070403 in both polarizations, thus making the FAST detection reliable. Considering the spectral baseline fluctuations appearing during the entire drift scan session, only ∼12 s observations around target transit time are utilized in our data reduction process. The final spectrum is an average of the ∼12 s transit data, and the background continuum is subtracted with similar smoothing and fitting method as the tracking data for sources described above.

The line centre of PGC 070403 determined by FAST, which is ∼7694 km s\(^{-1}\), is slightly redshifted compared with ALFALFA’s 7712 km s\(^{-1}\). Also, our line width is ∼16 per cent narrower than the Arecibo result, and the peak depth of the PGC 070403 absorbing feature obtained by FAST is more than 25 per cent shallower than ALFALFA’s value. Besides, the deepest structure in Arecibo data corresponding to a velocity of 7740 km s\(^{-1}\) does not show up in FAST observations. Combined with all of the factors mentioned above, the integrated optical depth obtained with FAST became ∼4.47 km s\(^{-1}\) smaller than ALFALFA’s 11.11 km s\(^{-1}\).

And as shown in the bottom panel of Fig. 10, with a four-channel bin applied to the original FAST data (which leads to a spectral resolution of ∼6 km s\(^{-1}\), comparing to ALFALFA’s ∼5.3 km s\(^{-1}\)), an absorbing feature with line centre at 7689.07 ± 2.14 km s\(^{-1}\), a narrower FWHM ≈ 64.37 ± 4.29 km s\(^{-1}\), a shallower \( S_{\text{HI,peak}} \) ≈ −8.525 ± 1.217 mJy, and a smaller \( \int \tau \text{d}v \approx 4.783 ± 0.809 \text{km s}^{-1} \), corresponding to an H I column density \( N_{\text{HI}} \approx (8.719 ± 1.570) \times 10^{20} \text{cm}^{-2} \), imposed on a background rms level of ∼1.21 mJy. Such a set of parameters exhibits noticeable differences with the ones acquired with the original, unbinned data.

As for the discrepancy of the low-frequency part of the line profiles between FAST and Arecibo observations, we suggest that polarization-dependent weak RFI or baseline fluctuations, which is quite common for FAST data (e.g. see fig. 26 in Jiang et al. 2020), could be a possible explanation. As can be seen in Fig. 9, compared with polarization A, the uncalibrated spectrum in polarization B...
shows a broader profile, as well as a greater depth in the low frequency regime, and is more similar to the ALFALFA result. When polarization-combined, the \( \sim 7740 \) km s\(^{-1}\) dip is largely missed due to polarization-dependent behaviours. Such an effect influences more for weak sources, especially when the integration time is limited. Thus, due to complications induced by high background noise level of drift scan observations, targeted long exposures with FAST or other sensitive radio telescopes should be needed to pin down the characteristics for PGC 070403. Anyway, current results still demonstrate the feasibility of identifying weak H\(_{\text{I}}\) absorption sources via FAST drift scan observations, thus more similar detections can be expected with the upcoming CRAFTS survey.

4 DISCUSSIONS

4.1 Note on ALFALFA detection rate

In our pilot study of extragalactic H\(_{\text{I}}\) absorbers with the FAST telescope during its commissioning phase, all of the five sources firstly discovered with 40 per cent of the ALFALFA data have been confirmed. Combining the five known sources identified by Wu et al. (2015) and Song et al. (in preparation), the detection rate for H\(_{\text{I}}\) absorption features in \( \alpha, 40 \), which is \( \sim 5.5 \) per cent, still holds, assuming the radio luminosity function of radio-loud AGNs to be in the form of that from Mauch & Sadler (2007). Such a rate is higher than the result estimated by Darling et al. (2011) with 7.4 per cent of ALFALFA data, which is \( \sim 3 \) per cent, as well as a later prediction by Allison et al. (2014) (\( \sim 1.6-4.4 \) per cent).

However, as noted in Saintonge (2007) and Haynes et al. (2011), the performance of the matched-filtering algorithm adopted by the source finding process drops significantly below an SNR value of \( \sim 6.5 \), which means that weak sources embedded in data cannot be identified completely, while Wu et al. (2015) detected all the 10 absorbers based on the same method described by Saintonge (2007). Since sources CGCG 049-033, PGC 070403, as well as three other known sources identified with \( \alpha, 40 \) exhibit flux levels less than the SNR threshold, the detection rate of 5.5 per cent can only be considered as a lower limit when estimating the total number of extragalactic H\(_{\text{I}}\) absorbers. However, since the rate for efficient detection (SNR > 6.5) is 2.75 per cent, which is only the half of the 5.5 per cent value, it is safe to say that at least 12 or 13 H\(_{\text{I}}\) absorbing systems can be identified with the \( \alpha, 100 \) data (Haynes et al. 2018), and the total number of detections is hard to predict due to algorithm limitations.

4.2 Implications on the prospect of H\(_{\text{I}}\) absorption line detections with FAST

Based on the different methods that we have tested during the pilot observations, we find that a single-pass scan can already resolve weak absorption, while the two-pass drift scan mode similar to that of ALFALFA (Giovanelli et al. 2005) would be a more efficient strategy to detect large number of H\(_{\text{I}}\) absorbers, since in this case, more time-varied RFI and other fluctuations can be excluded. In fact, the two-pass strategy is the observing mode already adopted by the FAST extragalactic H\(_{\text{I}}\) survey, which is one of the key projects undertaken at FAST, thus we can make good use of its large amount of data sets in the near future. Since it takes a single source \( \sim 12 \) s to pass through one beam during each scan, and a rotation angle of 23:4 would be applied to the 19-beam receiver (Li et al. 2018) to ensure maximum Dec. coverage and non-overlap between all feed horns during sky surveys, the total integration time for each source with all scans finished should be \( \sim 24 \) s. We calculate the mean rms value for a 24 s integration duration with all of the 19-beam data analysed by this paper, leading to a averaged noise level of \( \sim 3 \) mJy, which is comparable to the sensitivity of FLASH survey (\( \sim 3.2 \) mJy beam\(^{-1}\)) achieved with 2 hr of integration time by ASKAP (Allison et al. 2020). And such a noise level can be further suppressed to \( \sim 1.5 \) mJy with an extra four-channel bin along the frequency (velocity) direction, which is slightly lower than the noise level of WALLABY (\( \sim 1.6 \) mJy beam\(^{-1}\)) with 8 \times 2 hr ASKAP integration time, see Koribalski et al. (2020). Thus, in the following evaluations, we calculate the prospect of H\(_{\text{I}}\) absorption detections for FAST sky surveys with two sensitivity levels, \( \sim 3 \) and \( \sim 1.5 \) mJy, respectively. Take the averaged optical depth for 10 H\(_{\text{I}}\) absorbers (including five previously known samples) presented in Song et al. (in preparation), \( \tau \sim 0.352 \), as a typical value for extragalactic H\(_{\text{I}}\) absorption features, the corresponding normalized line depth should be of \( \sim 0.3 \) times the continuum flux. And suppose a reliable identification requires a \( \geq 5\sigma \) detection, averagely speaking, a continuum source, such as AGN, which served as the background for H\(_{\text{I}}\) absorption, is required to have a minimum flux of \( \sim 38 \) mJy for detecting with the original unsmoothed observations and \( \sim 25 \) mJy for the four-channel binned data.

The maximum zenith angle that can be attained by FAST is \( \sim 40^\circ \) (Nan et al. 2011; Li et al. 2018; Jiang et al. 2020). Given the telescope’s geographic latitude at 25° 39′ 10.6′ N, the observable Dec. extent of FAST should be \( \sim 14.35 < \delta < 65.65 \). Thus, the full FAST observable sky covers an area of \( \sim 23800 \) deg\(^2\), which is more than three times the mapped region of the ALFALFA survey (Giovanelli et al. 2005; Haynes et al. 2018). And the 19-beam receiver, the primary workhorse for sky surveys, operates at an H\(_{\text{I}}\) redshift range of \( \sim 6000 < z < 105 \) 000 km s\(^{-1}\), compared to ALFALFA’s \( \sim 2000 \) to 18 000 km s\(^{-1}\) (Haynes et al. 2018). Combined with both factors, the observable comoving volume of FAST should be \( \sim 4.2 \) Gpc\(^3\), which is nearly \( \sim 300 \) times the ALFALFA’s total coverage.

Following similar procedures by previous works such as Allison et al. (2014) and Wu et al. (2015), we adopt the local luminosity function for radio-loud AGNs proposed by Mauch & Sadler (2007). Ignoring the dependence of AGN distributions on redshift, we predict a total number of \( \sim 43000 \) AGNs with fluxes above the 5\( \sigma \) limit in the complete FAST sky survey coverage with a \( \sim 1.5 \) km s\(^{-1}\) spectral resolution, and \( \sim 49000 \) AGNs for the four-channel smoothed data. Applying the detection rate of H\(_{\text{I}}\) absorbing systems calculated by Song et al. (in preparation), \( \sim 2300 \) extragalactic H\(_{\text{I}}\) absorption lines should be identified with the complete set of original FAST sky survey data, and \( \sim 2600 \) such sources could be detected with the four-channel binned spectra. Of course, the calculations above only provide the most optimistic expectations, and the more or less compromised antenna gain of FAST for zenith angles larger than \( \sim 26.4^\circ \) (Li et al. 2018; Jiang et al. 2020) was not taken into account. Besides, a number of frequencies in the 1050–1450 MHz band are often heavily contaminated by RFI generated by distance measuring equipments for aviation, or emitters on board navigation satellites. Thus, the real number of H\(_{\text{I}}\) absorption line detection could be largely reduced. However, even neglecting the frequency band of \( \sim 1150–1300 \) MHz, which is most severely affected by RFI from satellites, \( \sim 1500 \) H\(_{\text{I}}\) absorption systems can still be expected from the unbinned data, and \( \sim 1700 \) from the four-channel binned data, which is consistent with the forecast made by Yu et al. (2017).

Another approach to predict the FAST detection prospect for H\(_{\text{I}}\) absorbers is with the NVSS (Condon et al. 1998) source count. Approximately \( 1.2 \times 10^5 \) continuum sources with \( S_{1.4\text{GHz}} > 38 \) mJy exist within the FAST observable sky; and the number for sources

\( S_{1.4\text{GHz}} > 38 \) mJy
with $S_{1.4GHz} > 25 \text{ mJy}$ is $\sim 1.8 \times 10^5$. With a detection rate of $\sim 5.5$ per cent applied, these numbers can lead to $\sim 6000$ extragalactic H\textsc{i} absorbers found within the original data, and $\sim 9000$ such systems for four-channel binned spectra. Still, since the NVSS catalogue does not provide associated redshift information for each individual source, and the chance of alignment between a high-$z$, H\textsc{i}-bearing galaxy is low, it is almost certain that a significant amount of H\textsc{i} absorbers associated with NVSS sources lie beyond the frequency coverage of FAST. Thus, such an estimation can only be considered as the rough upper limit for the CRAFTS survey.

5 CONCLUSIONS

In this paper, we report FAST observations of five extragalactic H\textsc{i} absorption lines first identified from the ALFALFA survey data. We confirm the existence of all absorption features from these sources. However, the line widths, optical depths, as well as H\textsc{i} column densities of the detected H\textsc{i} absorbers as derived from FAST data show noticeable discrepancies with the results previously obtained by Arecibo and WSRT, due to various factors. Since the FAST data have much higher SNR and finer spectral resolution compared with existing sky surveys, more features of the H\textsc{i} absorption lines, such as the double-horned structure of J1534+2513, can be revealed with high confidence. And the H\textsc{i} absorption line of PGC 070403, which exhibits the second-shallest line depth among the five, was successfully detected during a $\sim 12$ s integration time using the drift scan mode, which will be the final choice for the upcoming extragalactic H\textsc{i} surveys. These observations, which can be considered as the first batch of extragalactic H\textsc{i} absorption line detections performed by FAST, demonstrated the capability of this telescope for H\textsc{i} absorption studies. It is expected that with a larger sky coverage and higher sensitivity than Arecibo, over 1500 extragalactic H\textsc{i} absorbers could be unveiled with the entire set of FAST extragalactic sky survey data.

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DATA AVAILABILITY

The data analysed by this work are from the FAST project nos 3017 (UGC 00613, ASK 378291.0, CGCG 049-033, J1534+2513) and N2020.3 (PGC 070403), and can be accessed by sending request to the FAST Data Centre or to the corresponding authors of this paper.

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